

Probing EWIMPs with Drell–Yan process at 100 TeV colliders

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Based on **1810.07349** with S. Chigusa and T. Moroi (Tokyo U.)



Introduction

ElectroWeakly Interacting Massive Particle (EWIMP)

- EWIMP: massive particles charged under $SU(2)_L$.
- Motivated by many Standard Model extensions.
- Good candidate for Dark Matter.

➔ Extensive searches being performed.

MSSM

- Higgsino: $SU(2)$ doublet

$m \simeq 1 \text{ TeV}$ for thermal relic.

e.g. "natural SUSY" [Baer+ 12; ...]

- Wino: $SU(2)$ triplet

$m \simeq 3 \text{ TeV}$ for thermal relic.

e.g. "mini-split" [Arvanitaki+ 12; ...]

Minimal DM

[Cirelli+ 05, 07]

- 5-plet fermion

$m \simeq 10 \text{ TeV}$ for thermal relic.

- 7-plet scalar

$m \simeq 25 \text{ TeV}$ for thermal relic.

* additional stabilization for 7-plet [Del Nobile+ 15]

Conventional search

DM detection

- direct detection

wino, MDM: well within future prospect, higgsino: under neutrino floor

- indirect detection

wino, MDM: strongly constrained

subject to astrophysical uncertainties

Collider search

- disappearing charged track

wino: strong way to probe

depends strongly on lifetime, or mass difference

- mono-X search

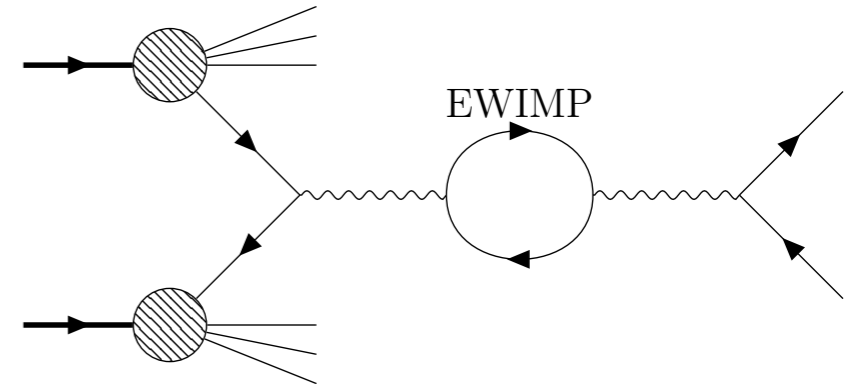
* Higgsino is basically more difficult to probe.

Collider “indirect” search

- Today’s topic: Drell-Yan process @ 100 TeV hadron collider.

* Focus on $ee, \mu\mu$ channels in this talk.

[See also Harigaya+ 15; Matsumoto+ 17, 18; Di Luzio+ 18]



- Not directly produce EWIMPs = “indirect” search.
 - ✓ insensitive to lifetime and decay products.
 - ✗ EWIMP effects only @ NLO
 - ✗ Have to deal with backgrounds and systematics.
- Offers complementary information, to say the least.

Outline

1. Introduction

2. EWIMP effect on Drell-Yan

3. Detection reach

4. Summary

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1. Introduction

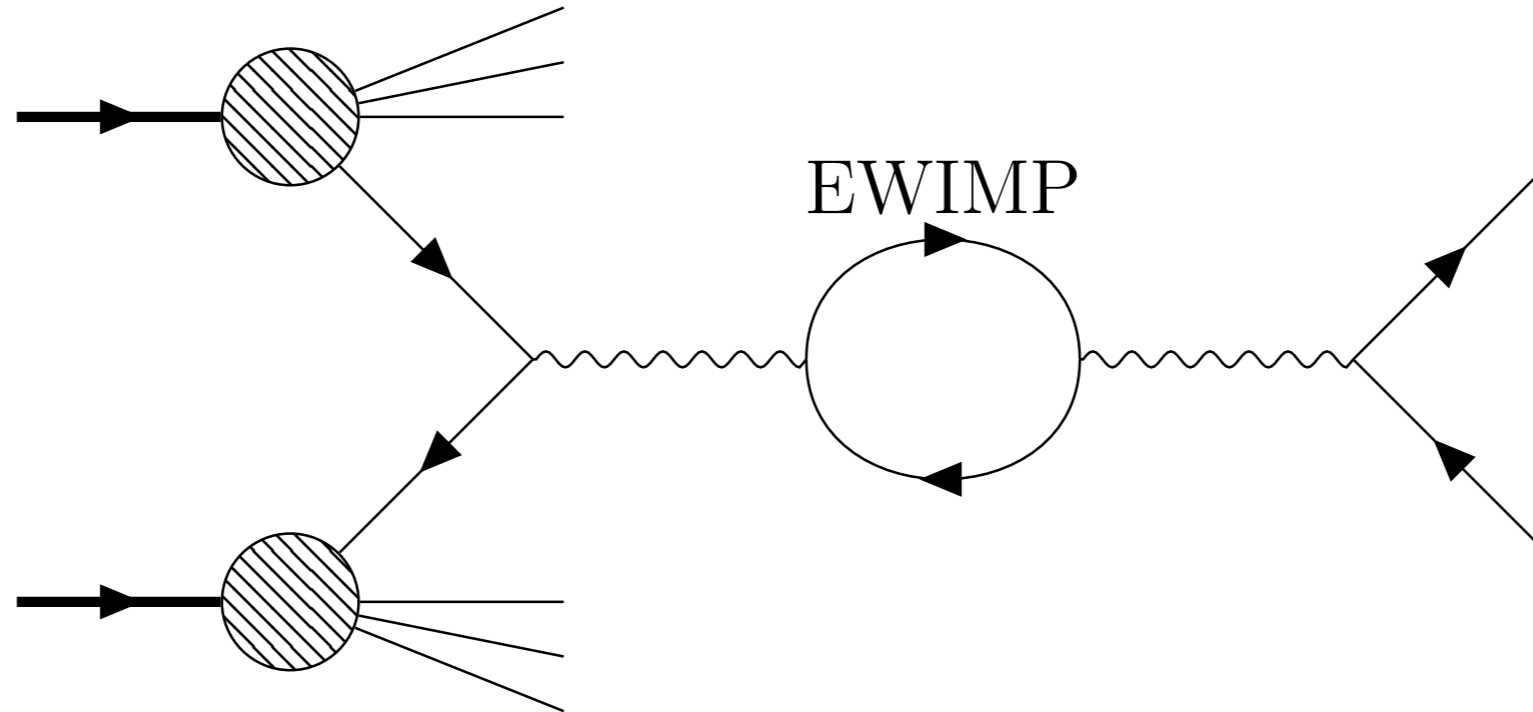
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Oblique correction

- EWIMP acts as an oblique correction:



- We take into account the full contribution.

← important to distinguish from systematics (discussed later).

c.f. $S, T, U, (W, Y, \dots)$

for $\sqrt{s} \ll m_{\text{NP}}$,

[Peskin&Takeuchi 90,92; Barbieri+ 04; ...]

and the gauge coupling running for $\sqrt{s} \gg m_{\text{NP}}$.

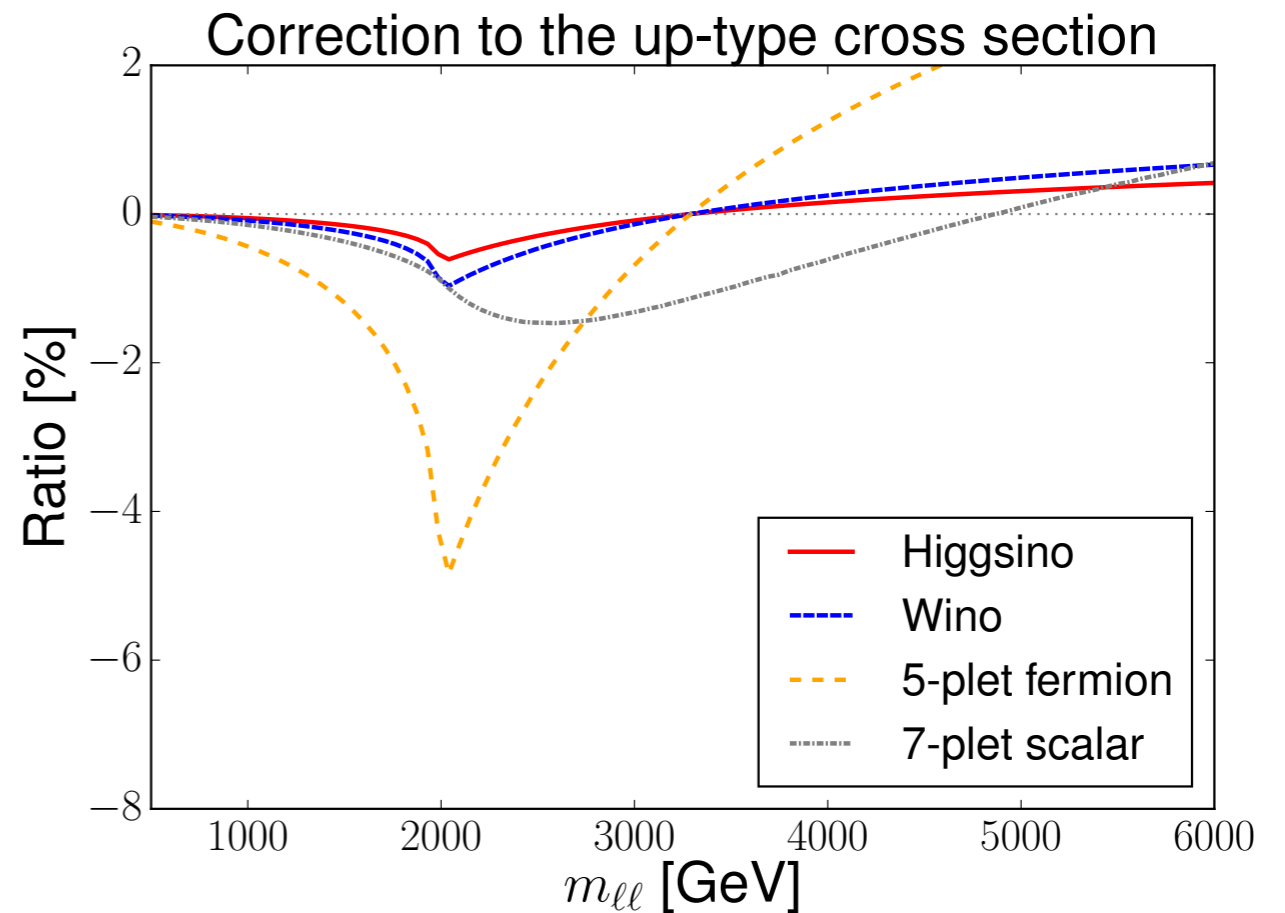
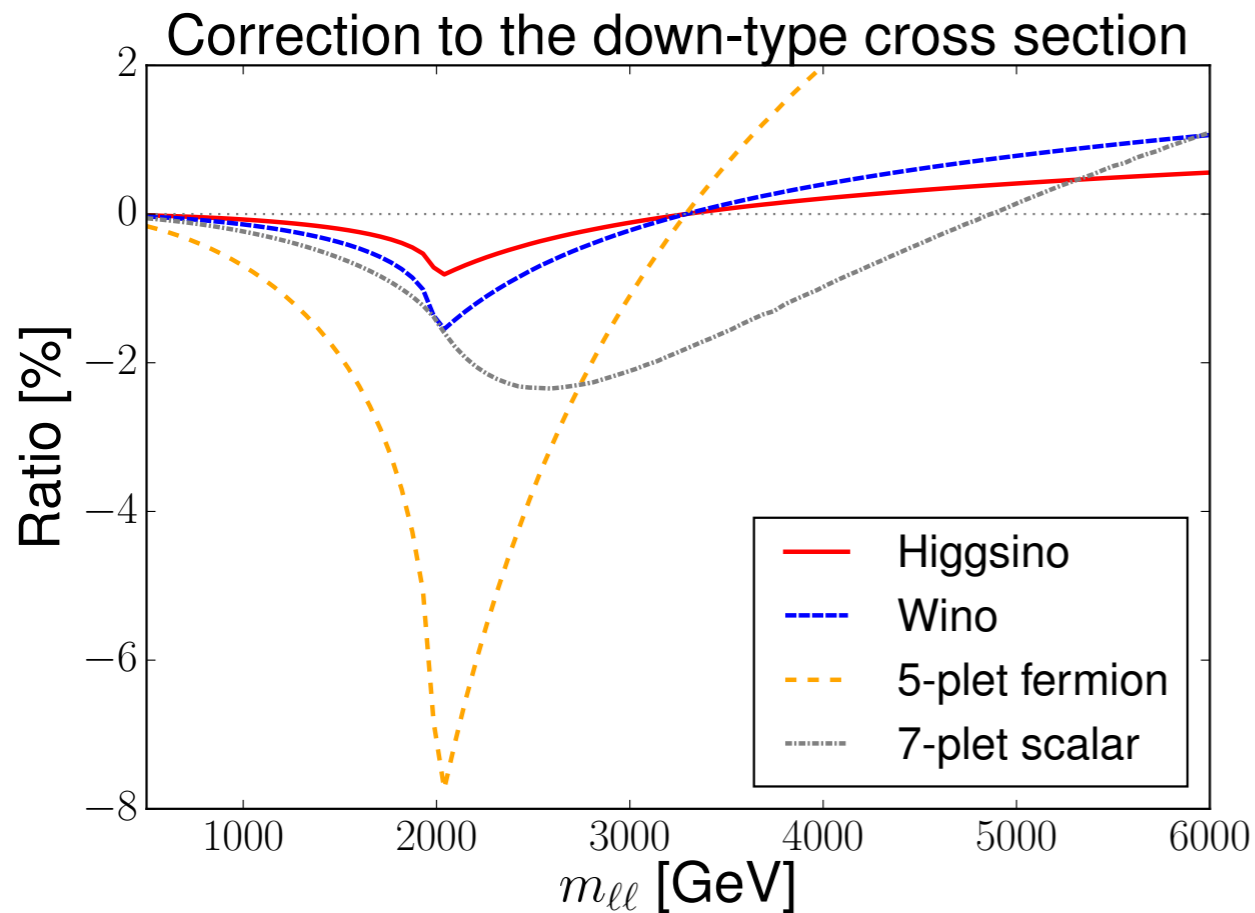
[Alves+ 14; Gross+ 16; ...]

EWIMP effect

- Define the EWIMP correction

$$\delta_\sigma(m_{ll}) \equiv \frac{d\sigma_{\text{EWIMP}}/dm_{ll}}{d\sigma_{\text{SM}}/dm_{ll}} \quad \text{for each parton.}$$

- Plot of $\delta_\sigma(m_{ll})$ for $m_{\text{EWIMP}} = 1 \text{ TeV}$:



- Correction peaked at $m_{ll}/2 = m_{\text{EWIMP}} \rightarrow$ suitable for 100 TeV collider!

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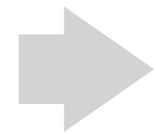
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Event generation and model

- Experiment not performed yet.



Events generated by Monte-Carlo (MadGraph5 & Pythia8 & Delphes3)

$$\mathcal{L} = 30 \text{ ab}^{-1}, \quad 500 \text{ GeV} < m_{ll} < 15 \text{ TeV}, \quad 145 \text{ bin}$$

* NLO: event # $\rightarrow \times 2$, detector: $\times 0.8$ and smear momentum.

- Theoretical prediction also computed by MC.

$$\text{Event \#}: \tilde{x}_i(\mu) = \sum_{m_{ll}^{(i)} < m_{ll}^{\text{obs}} < m_{ll}^{(i+1)}} [1 + \mu \delta_\sigma(m_{ll}^{\text{true}})] \quad \text{for } i\text{-th bin,}$$

where m_{ll}^{obs} : observed invariant mass (simulated by Delphes),

and m_{ll}^{true} : true invariant mass (read from MadGraph output).

* Pure SM limit: $\mu = 0$, full EWIMP effect: $\mu = 1$.

Fitting based approach

- Many sources for systematic uncertainties:

luminosity, beam energy, PDF choice, factorization/renormalization scale, ...

➔ Introduce nuisance parameters $\theta = \{\theta_i\}$ to absorb them:

$$\tilde{x}_i(\mu, \theta) = \tilde{x}_i(\mu) f_{\text{sys},i}(\theta), \quad f_{\text{sys},i}(0) = 1.$$

- Variation of nuisance parameters = systematics uncertainties.
- For illustration, we take in the following

$$f_{\text{sys},i} = e^{\theta_1} (1 + \theta_2 p_i) p_i^{(\theta_3 + \theta_4 \ln p_i + \theta_5 \ln^2 p_i)}, \quad p_i = 2m_{ll,i}/\sqrt{s},$$

that is known to work for LHC. [CMS collaboration 08]

* This choice is just an example.

Fitting based approach

- Define test statistics as

[Cowan+ 10]

$$q_0 \equiv -2 \ln \frac{L(\mathbf{x}; \mu = 0, \hat{\hat{\theta}})}{L(\mathbf{x}; \hat{\mu}, \hat{\theta})},$$

where

$$L(\mathbf{x}; \mu, \theta) \equiv \underbrace{\left(\prod_i \exp \left[-\frac{(x_i - \tilde{x}_i(\mu, \theta))^2}{2\tilde{x}_i(\mu, \theta)} \right] \right)}_{\text{likelihood in signal region}} \underbrace{\left(\prod_{\alpha} \exp \left[-\frac{\theta_{\alpha}^2}{2\sigma_{\alpha}^2} \right] \right)}_{\text{our knowledge on systematics}},$$

$\hat{\hat{\theta}}$: maximize $L(\mathbf{x}; \mu = 0, \theta)$, $(\hat{\mu}, \hat{\theta})$: maximize $L(\mathbf{x}; \mu, \theta)$.

$\sqrt{q_0} = n$ corresponds to $n\sigma$ discovery. [Wilk 38]

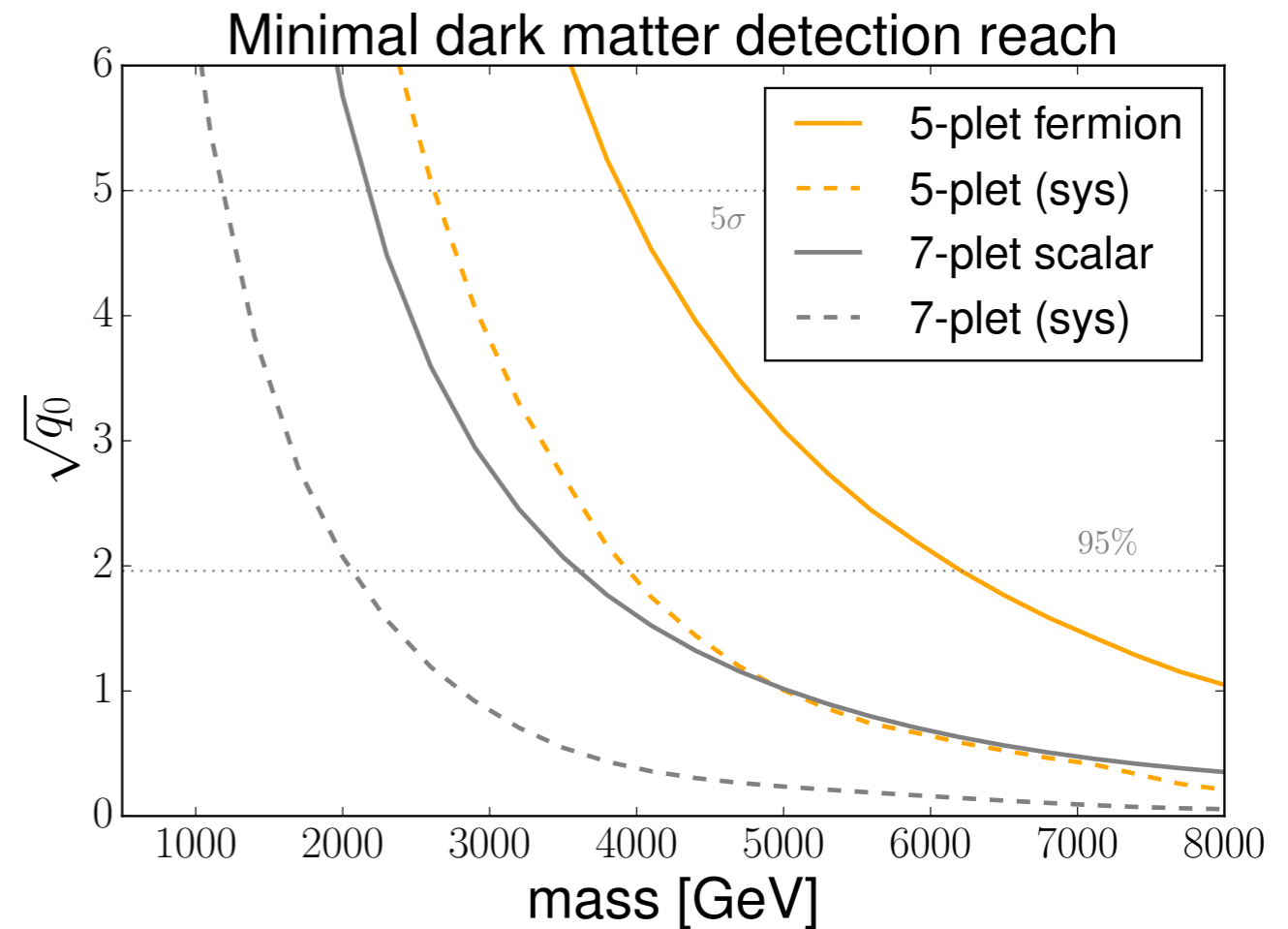
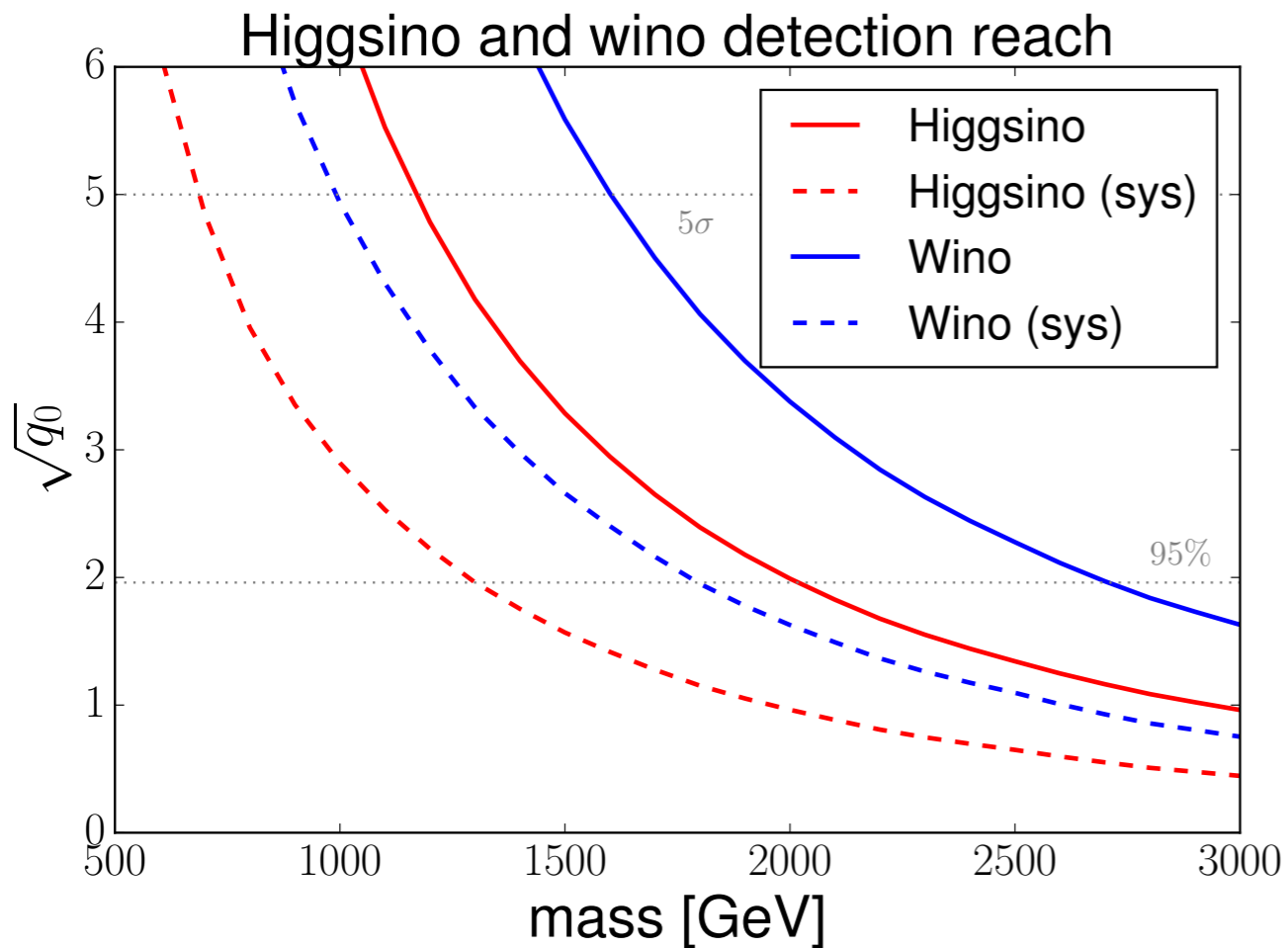
- Estimated the variance σ by varying one of the followings:

- luminosity: $\pm 5\%$
- beam energy: $\pm 1\%$
- renormalization scale: $2Q, Q/2$
- factorization scale: $2Q, Q/2$
- PDF choice

* They are well fitted by our choice of $f_{\text{sys},i}$.

Result

$$\mathcal{L} = 30 \text{ ab}^{-1}$$



	Higgsino	Wino	5-plet	7-plet
5 σ	0.7 TeV	1.0 TeV	2.6 TeV	1.2 TeV
95% C.L.	1.3 TeV	1.8 TeV	3.9 TeV	2.1 TeV

Comparison

$$\mathcal{L} = 30 \text{ ab}^{-1}$$

- **Higgsino:**

	mono-jet	disappearing track	indirect
5σ	0.2 - 0.5 TeV	0.9 - 1.3 TeV	0.7 (- 1.2) TeV
95% C.L.	0.9 - 1.4 TeV	1.1 - 1.5 TeV	1.3 (- 2.0) TeV

Indirect method can be strongest!

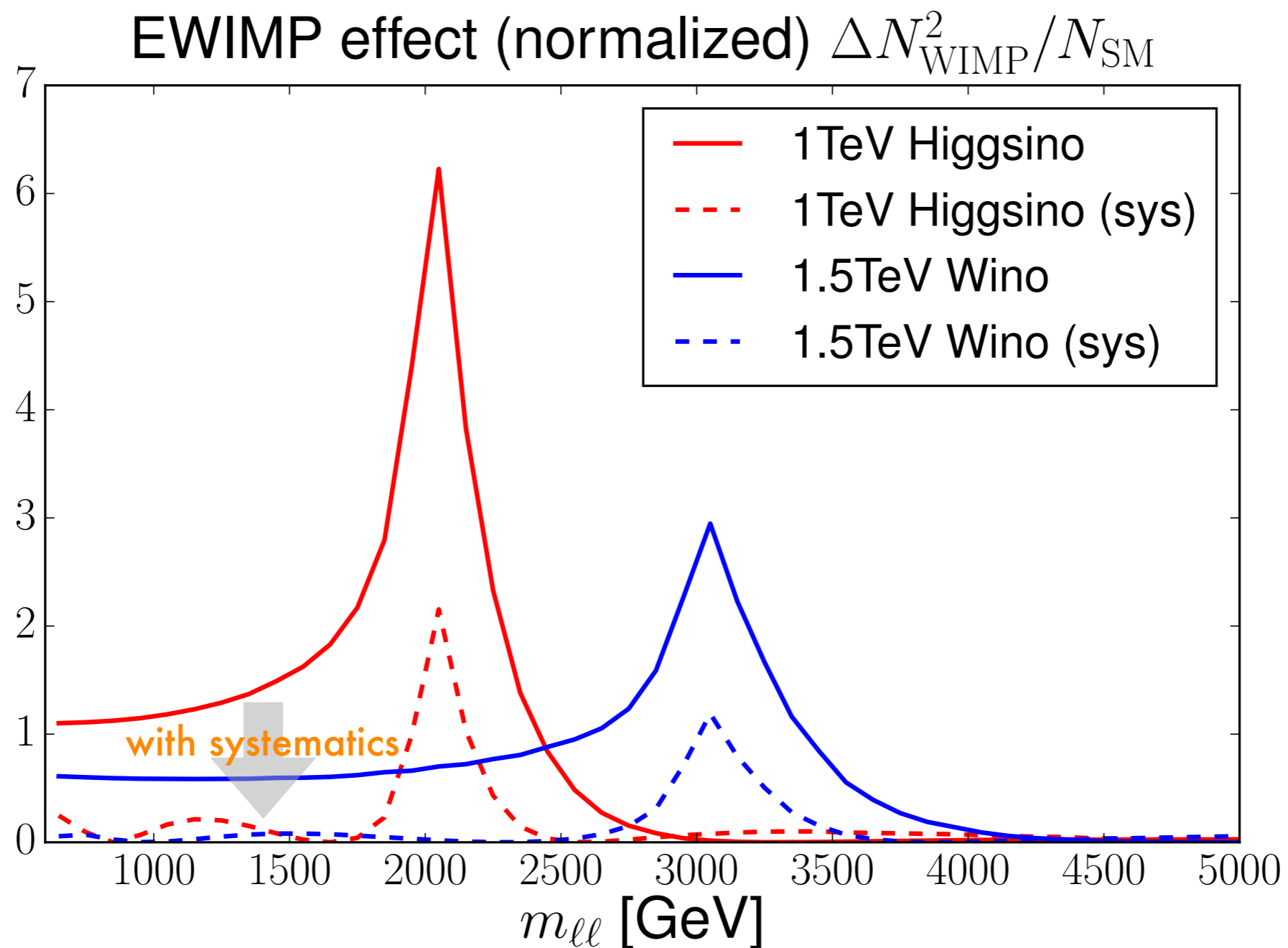
- **Wino:**

	mono-jet	disappearing track	indirect
5σ	0.5 - 1.0 TeV	3.6 - 5.4 TeV	1.0 (- 1.6) TeV
95% C.L.	1.5 - 2.0 TeV	4.5 - 6.6 TeV	1.8 (- 2.7) TeV

**Disappearing track seems best way,
indirect method still provides complementary information.**

* Mono-jet and disappearing track taken from 1805.00015.

Contribution from each bin



- Peak is the dominant contribution.
- Peak is not absorbed by systematics.

Outline

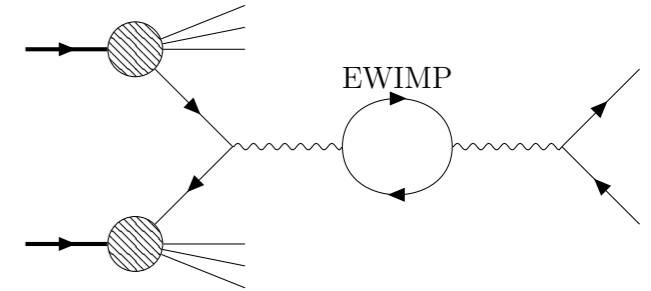
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Summary



- Introduced “indirect” method to probe EWIMPs.
- Study prospect at future 100 TeV hadron collider.
- Indirect method can be the strongest probe for Higgsino.
- Provides complementary information for other EWIMPs.
- Reducing systematic uncertainties has a significant effect.

	Higgsino	Wino	5-plet	7-plet
5σ	0.7 TeV	1.0 TeV	2.6 TeV	1.2 TeV
5σ (opt.)	1.2 TeV	1.6 TeV	3.9 TeV	2.2 TeV
95% C.L.	1.3 TeV	1.8 TeV	3.9 TeV	2.1 TeV
95% C.L. (opt.)	2.0 TeV	2.7 TeV	6.2 TeV	3.6 TeV

Back up

Conventional DM search

- **DM direct detection**

[Hil and Solon 13; Hisano+ 15; XENON collaboration 18; ...]

- ✓ **Wino, MDM: well within the future detection reach.**
- ✗ **Higgsino: below the neutrino floor, difficult to detect.**
(assuming mass splitting (pseudo Dirac) to avoid Z-boson exchange)

- **DM indirect detection**

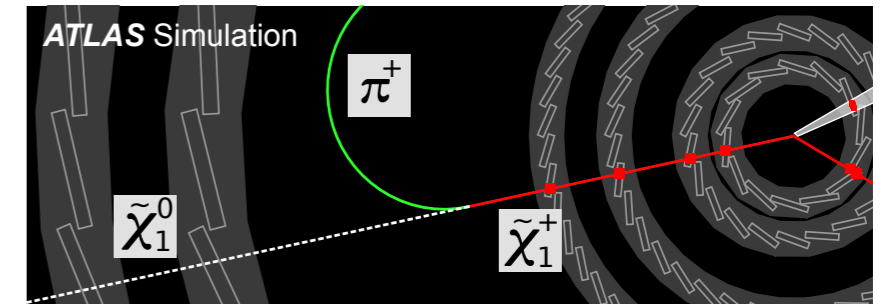
[Fermi-LAT collaboration 15; HESS collaboration 16,18; ...]

- ✓ **Wino, MDM: severely constrained.** (or even excluded for cuspy DM profile)
- ✗ **Higgsino: difficult to probe due to smaller coupling.**
- ✗ **Uncertainties of astrophysical origin, e.g. DM profile.**

* EWIMP must be the dominant component of DM for both methods.

Conventional collider search

- Disappearing charged track:



- ✓ For pure wino/higgsino: $m_{\tilde{W}} < 460 \text{ GeV}$, $m_{\tilde{h}} < 152 \text{ GeV}$.

[ATLAS 17; CMS 18]

- ✓ Wino thermal relic can be well covered by future colliders.

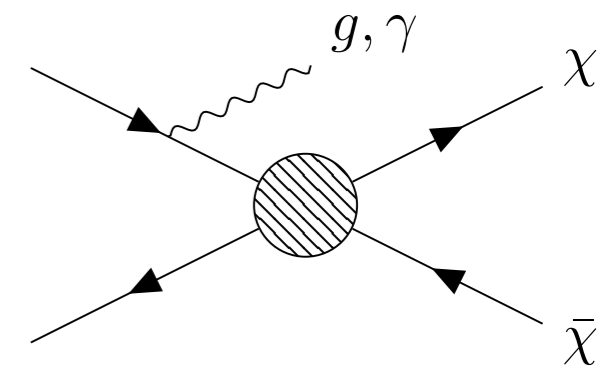
- ✗ strongly depends on lifetime, or mass difference.

* higgsino: dim-5, and wino: dim-7 in terms of EFT.

- Mono-X search:

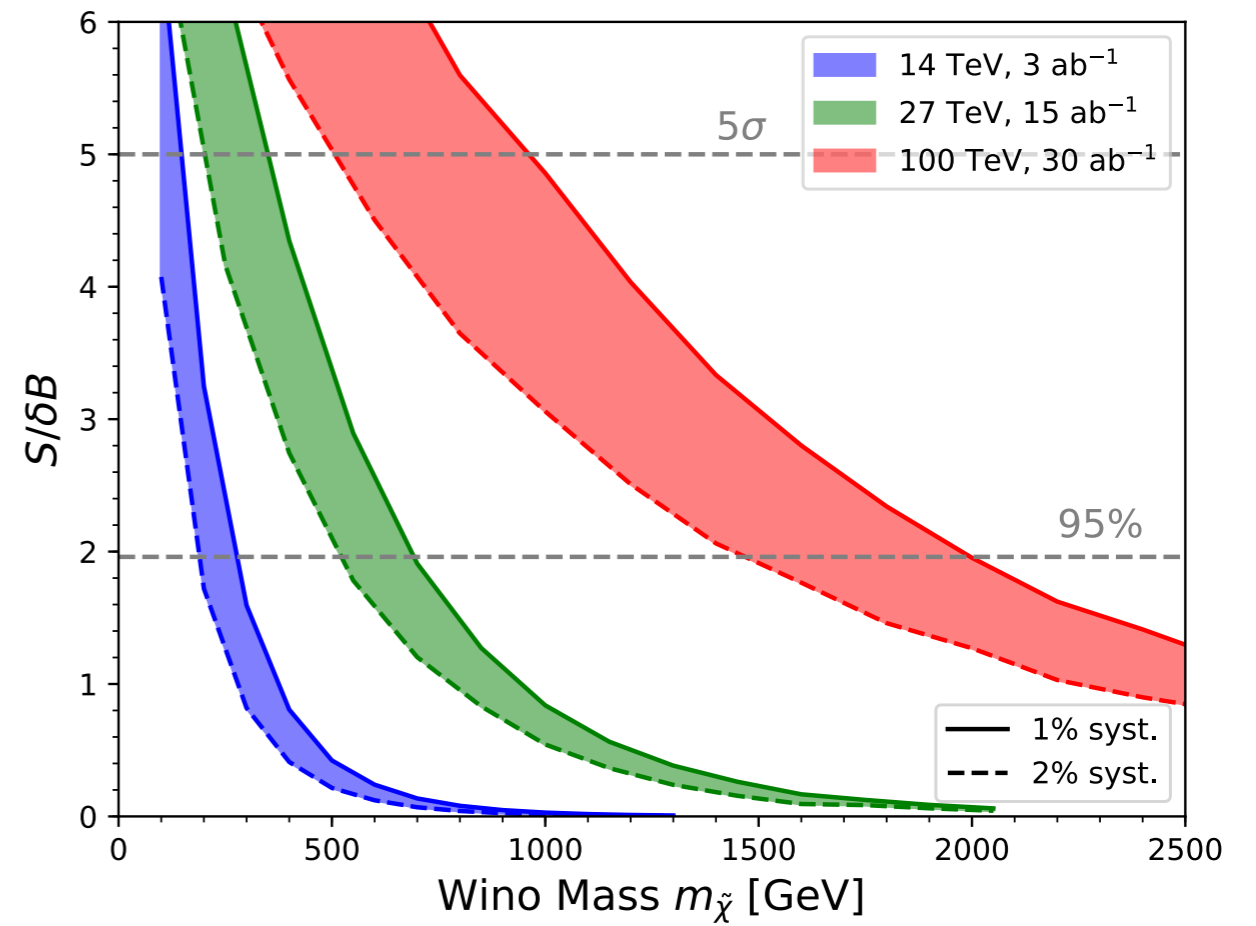
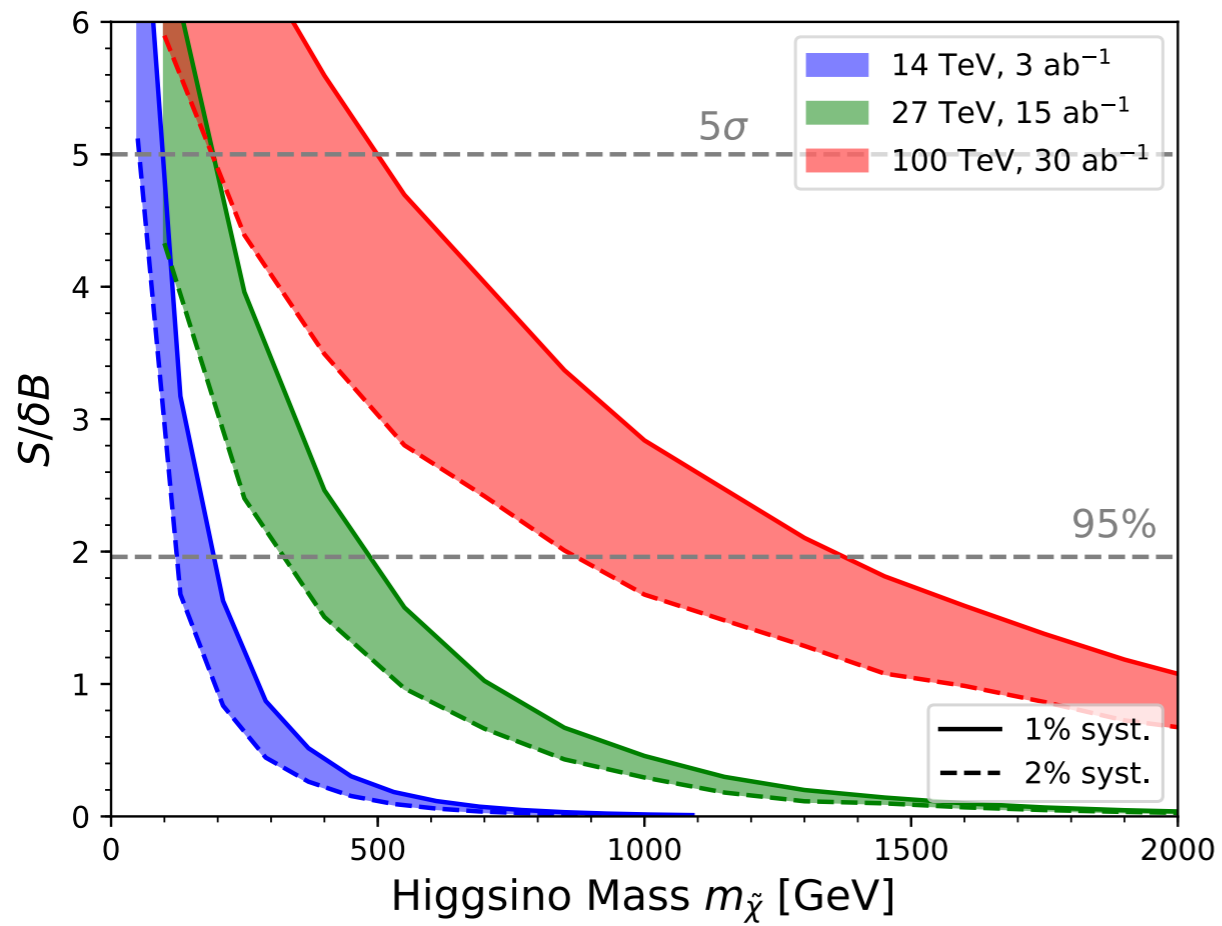
- ✗ Signal and background: similar distribution

e.g. no significant bound on higgsino at LHC14 [Baer+ 14; ...]



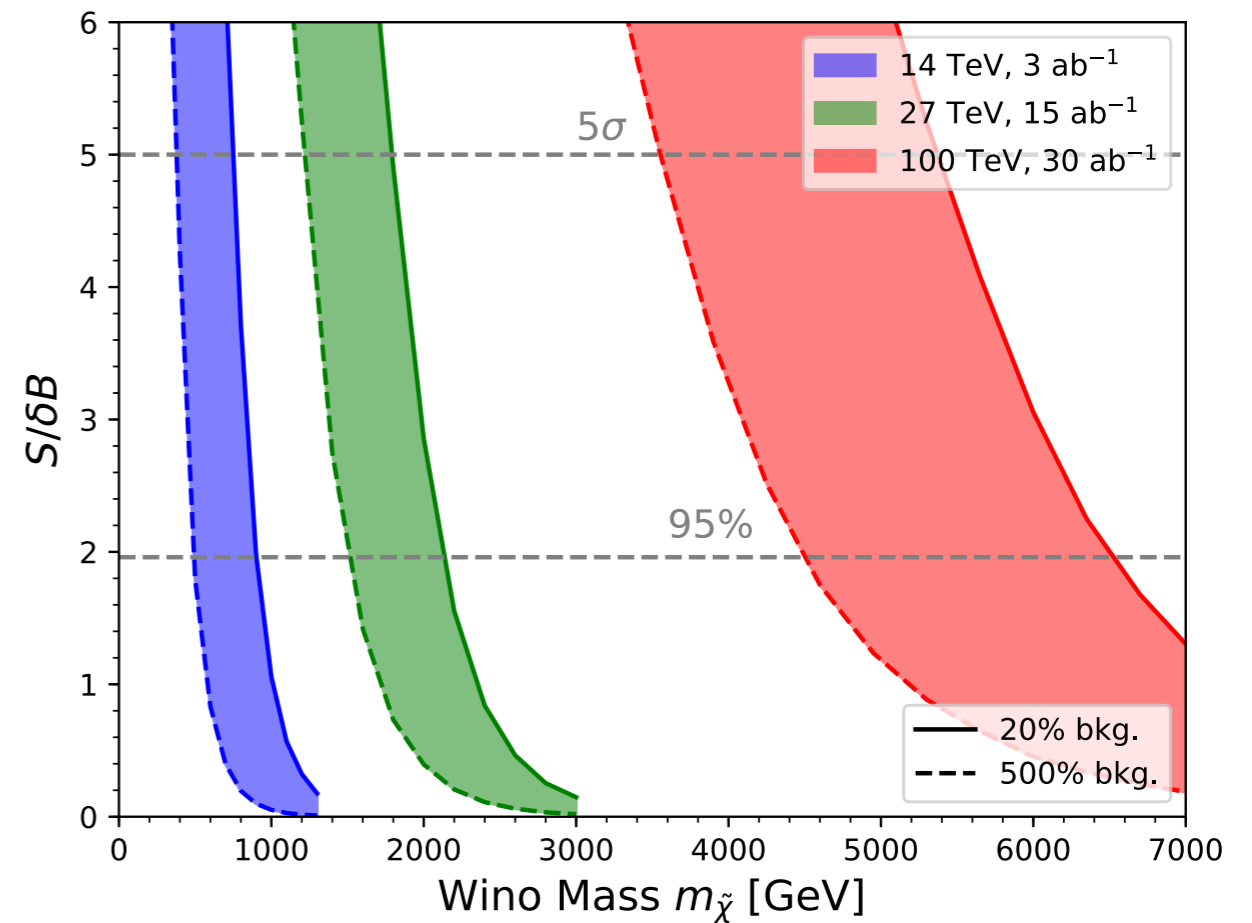
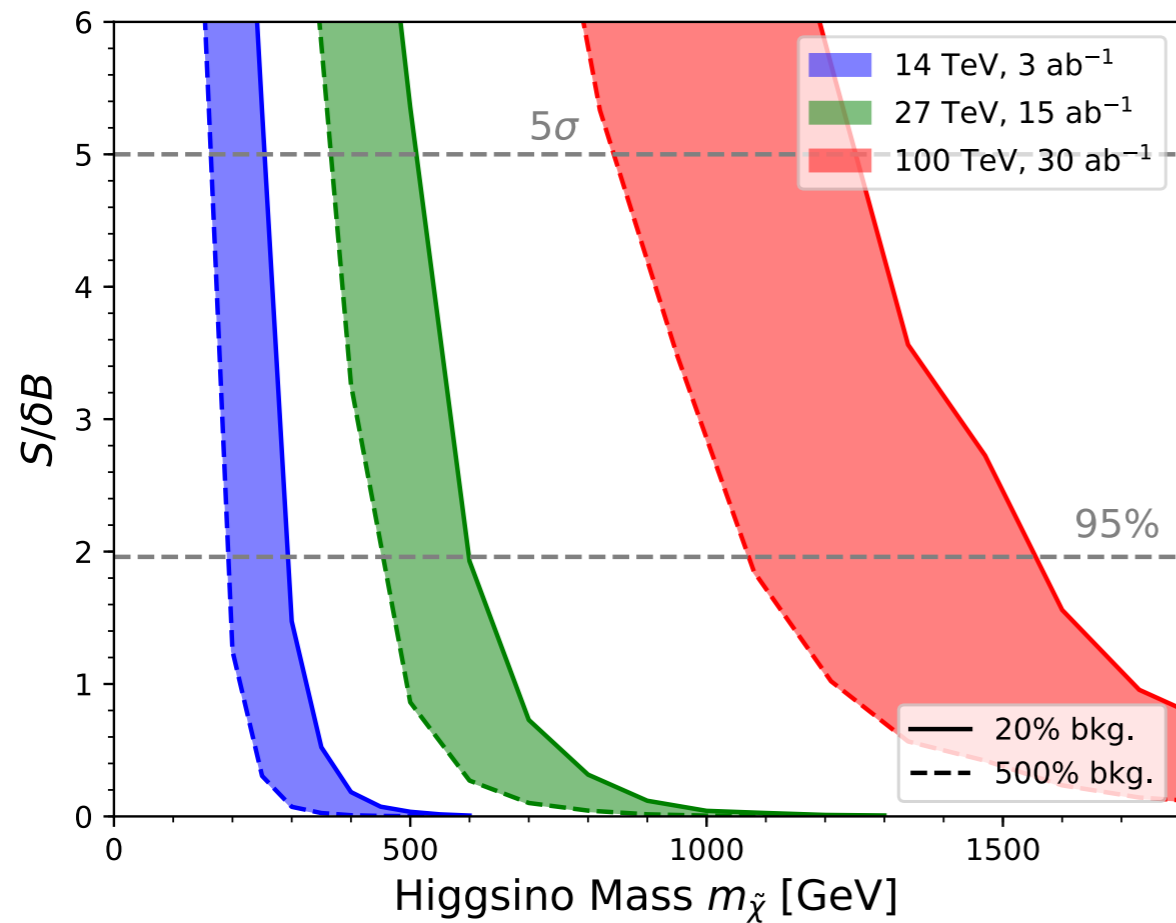
Collider: mono-jet

[Han+ 18]



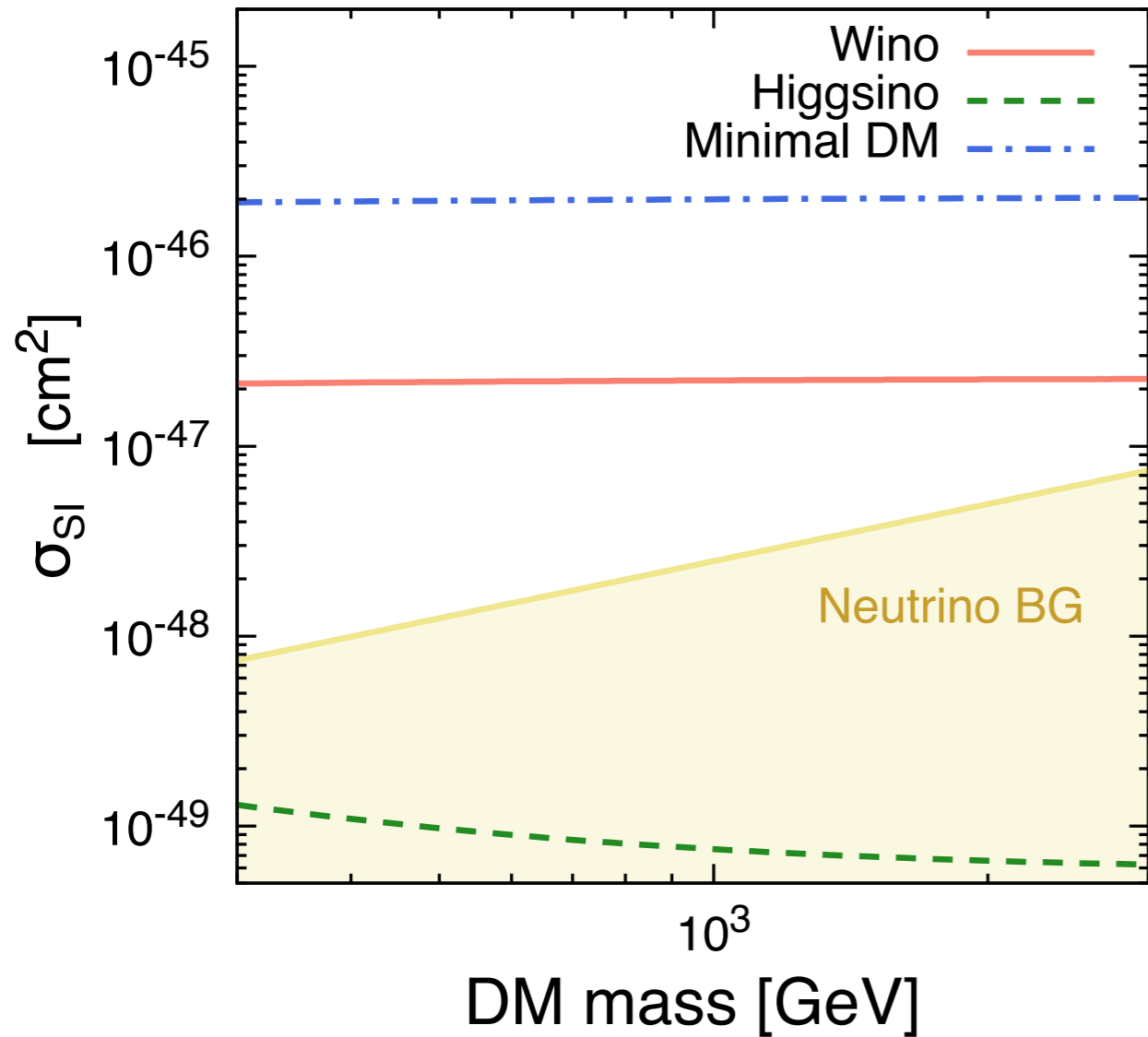
Collider: disappearing track

[Han+ 18]

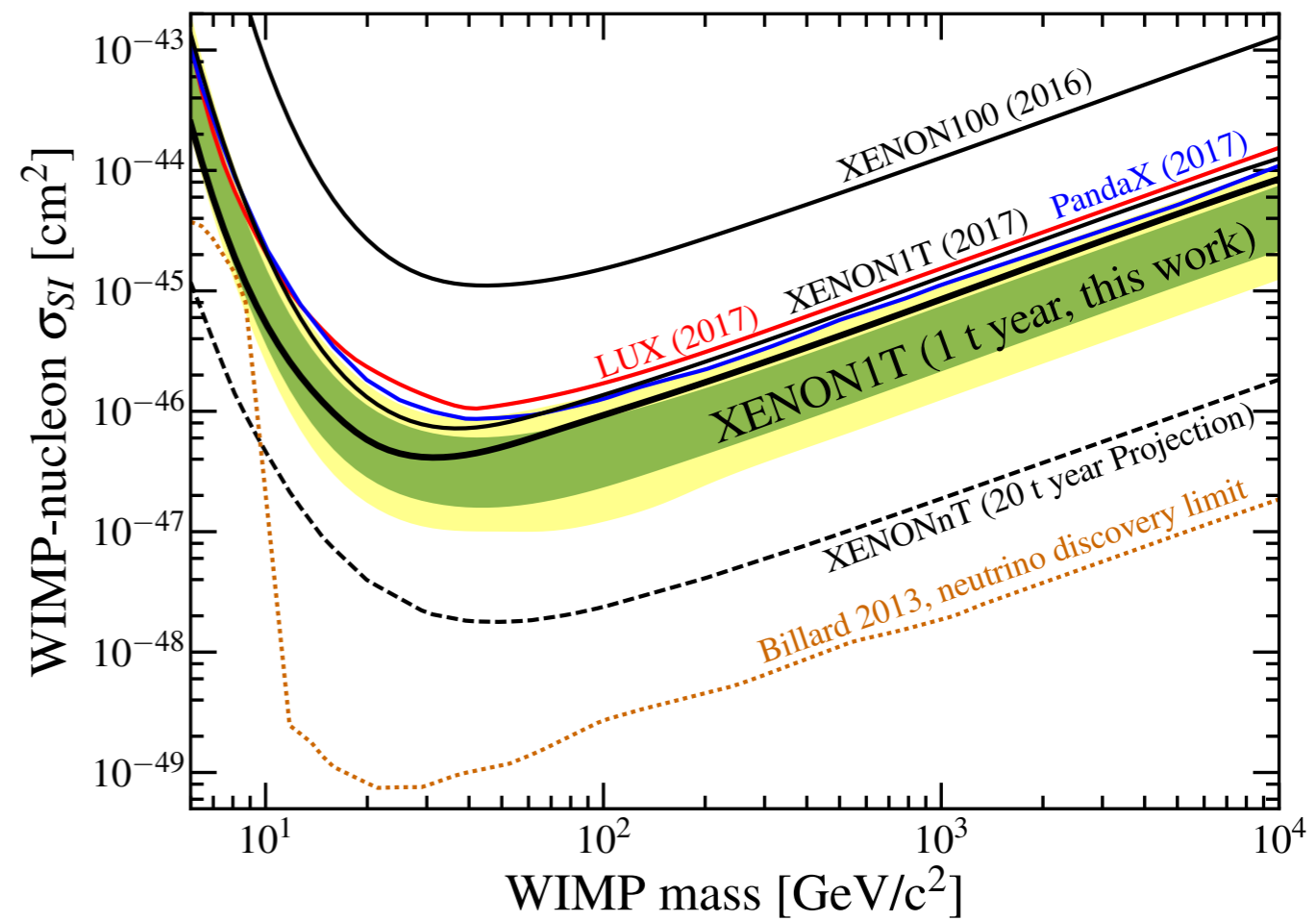


DM direct detection

[Hisano+ 15]

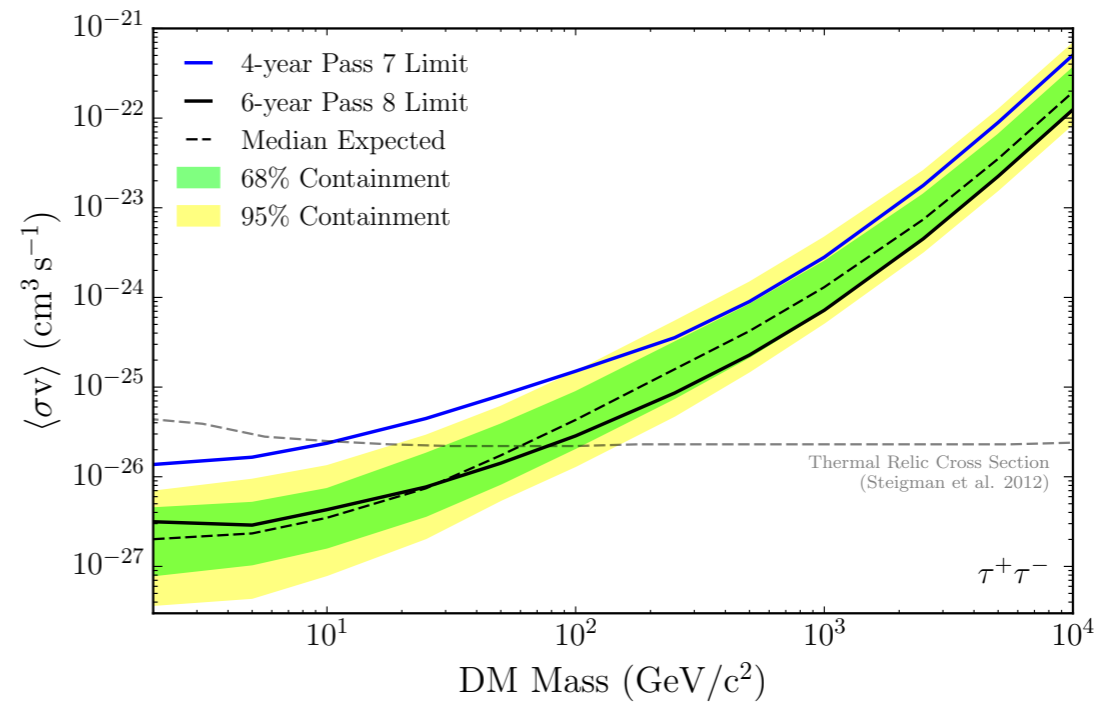
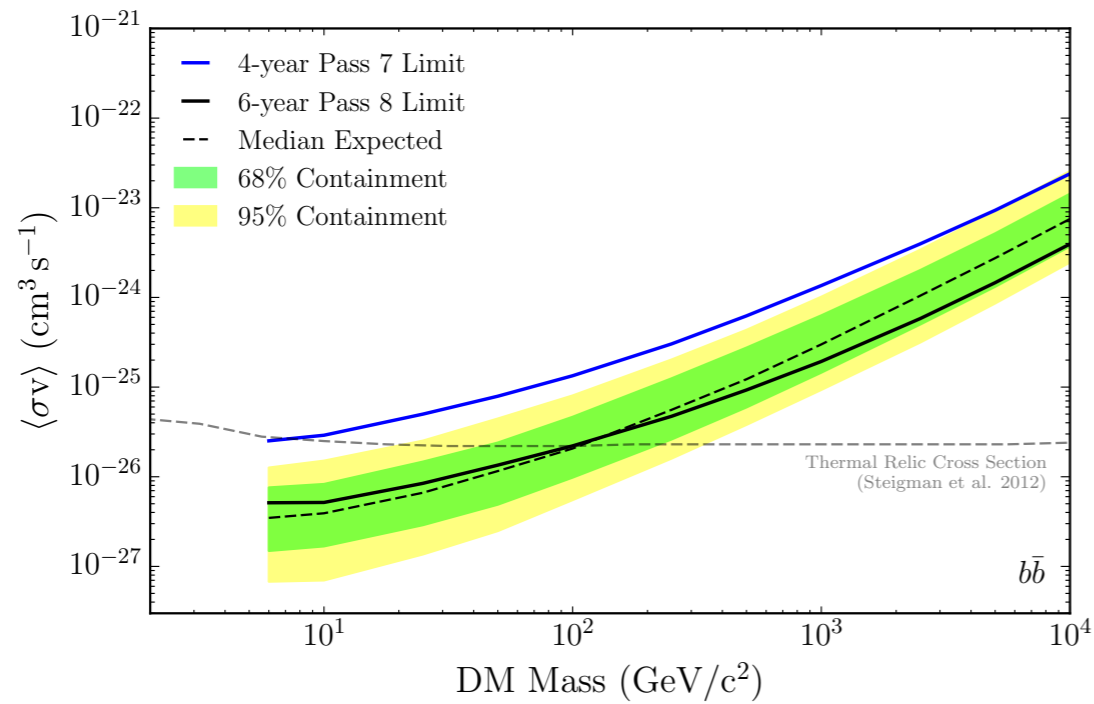


[XENON collaboration 18]



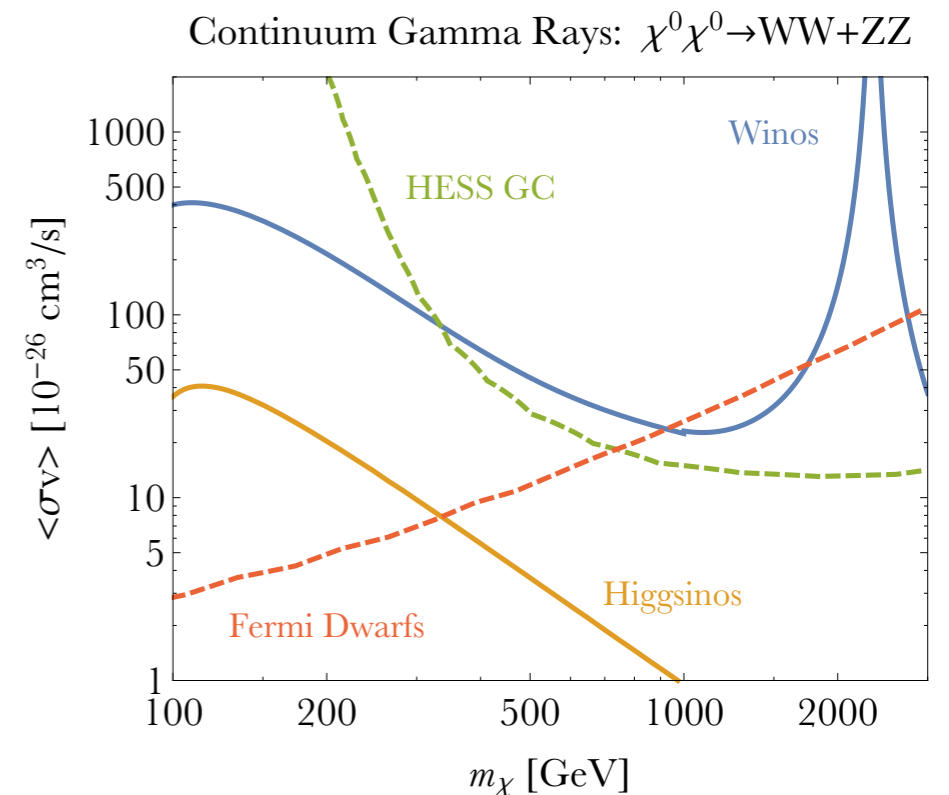
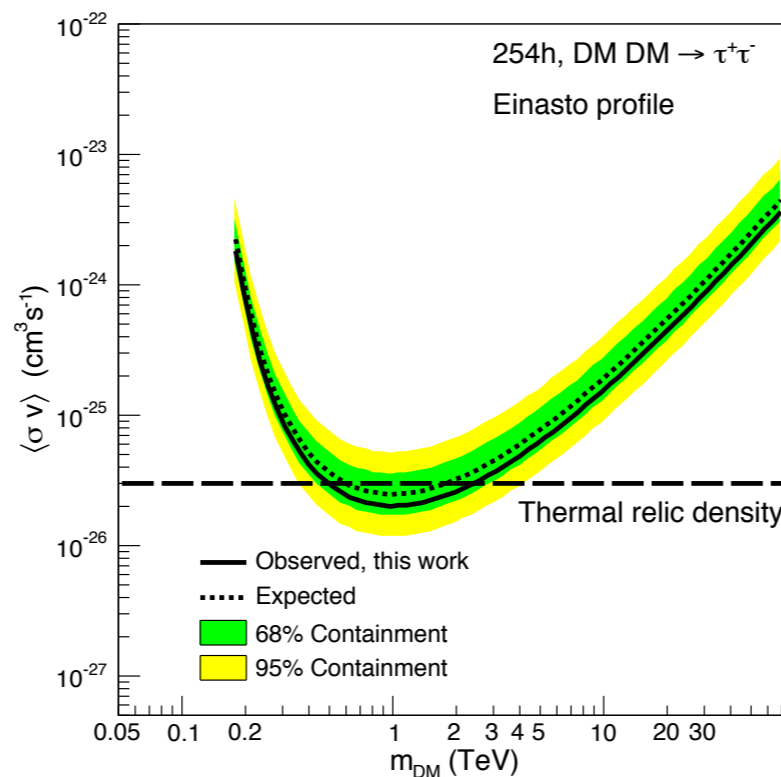
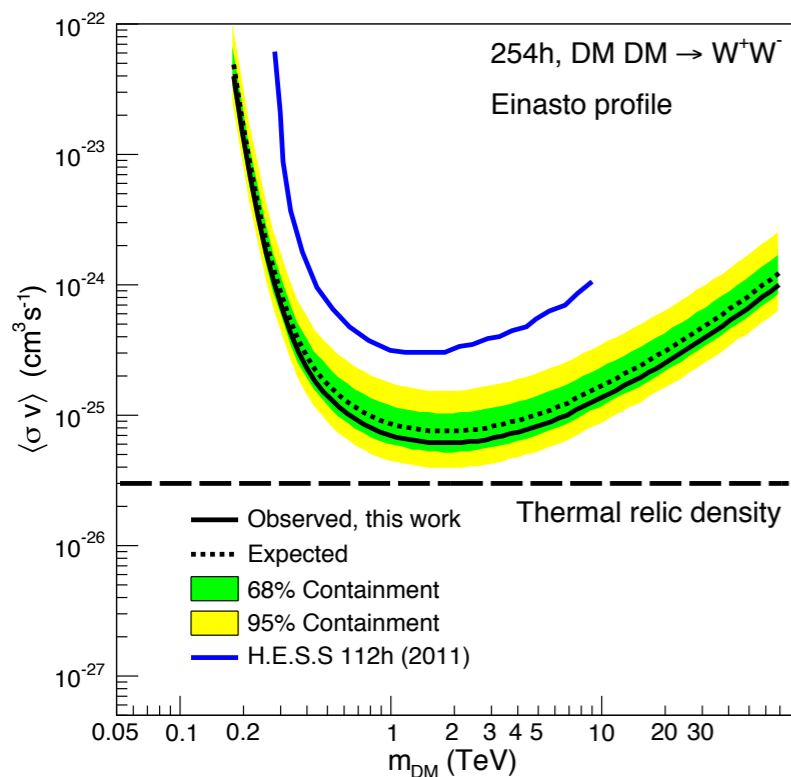
DM indirect detection

dwarf spheroidal galaxy [Fermi-LAT collaboration 15]



galactic center [HESS collaboration 16]

[Krall and Reece 17]



* Einasto (cuspy) profile is assumed for HESS.

Mass determination

